

**Partial Differential Equations – Math 442 C13/C14**  
**Fall 2009**  
**Homework 4 — due October 9**

1. Consider the boundary value problem

$$A'' + \lambda A = 0, \quad A'(0) + aA(0) = 0, \quad A(L) = 0.$$

- (a) Show that if  $a < 0$ , then there is no negative eigenvalue.
- (b) Under which conditions is there a zero eigenvalue?
- (c) Show there are infinitely many positive eigenvalues for any value of  $a$ .

**Bonus:** We showed in (a) that if  $a < 0$  then there is no negative eigenvalue. It turns out that for some positive  $a$ , this problem has a negative eigenvalue (and for some others it does not). Write down a condition on  $a$  which determines whether such an eigenvalue exists.

*Solution.*

- (a) Let's say we have a negative eigenvalue  $\lambda < 0$ . Then if we say  $A(x) = e^{rx}$ , we have

$$r^2 + \lambda = 0, \\ r = \pm\sqrt{-\lambda}.$$

Since  $\lambda < 0$ , this means the roots are real, let us write them as  $r = \pm b$  where  $b > 0$ . Then we have

$$A(x) = C_1 e^{bx} + C_2 e^{-bx}.$$

Plugging in the initial conditions gives

$$C_1 b - C_2 b + a(C_1 + C_2) = 0, \\ C_1 e^{bL} + C_2 e^{-bL} = 0.$$

The second equation can be written as  $C_2 = -C_1 e^{2bL}$ , and plugging this into the first gives

$$C_1(a + b) + C_1 e^{2bL}(b - a) = C_1(b(1 + e^{2bL}) + a(1 - e^{2bL})).$$

This must be zero, but  $C_1$  cannot be zero (since then so would  $C_2$  be), so we need to solve

$$b(1 + e^{2bL}) + a(1 - e^{2bL}) = 0. \tag{1}$$

However, this equation has no solution! Notice that since  $bL > 0$ , we know  $e^{2bL} > 1$ , and thus  $1 - e^{2bL} < 0$ . Since  $a$  is also less than zero, we know that

$$a(1 - e^{2bL}) > 0.$$

But since  $b(1 + e^{2bL}) > 0$  as well, it is not possible to solve (1).

- (b) If we have a zero eigenvalue, this means we have a solution to  $A''(x) = 0$  with those boundary conditions. However, this means that  $A(x) = \alpha x + \beta$ , and plugging this into the boundary conditions gives

$$\alpha + a\beta = 0, \\ \alpha L + \beta = 0.$$

We want to find a nontrivial solution to this system, and if we write it as a matrix equation:

$$\begin{pmatrix} 1 & a \\ L & 1 \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix},$$

and this has a nontrivial solution iff the matrix is singular, i.e. if its determinant is zero. So the condition is that

$$\det \begin{pmatrix} 1 & a \\ L & 1 \end{pmatrix} = 1 - aL = 0,$$

or  $aL = 1$ .

(c) If  $\lambda > 0$ , then the solutions to the ODE look like

$$C_1 \cos(\omega x) + C_2 \sin(\omega x),$$

where  $\omega^2 = \lambda$ . Plugging in the boundary conditions gives

$$\begin{aligned} C_2 \omega + aC_1 &= 0, \\ C_1 \cos(\omega L) + C_2 \sin(\omega L) &= 0. \end{aligned}$$

Solving the first equation gives  $C_2 = -aC_1/\omega$ , and plugging this into the second and doing some algebra gives

$$\tan(\omega L) = \frac{\omega}{a}.$$

The question is, how many roots does this equation have? We cannot answer this analytically, but we can see from the graph that it has to have infinitely many; see an example picture in Figure 1.

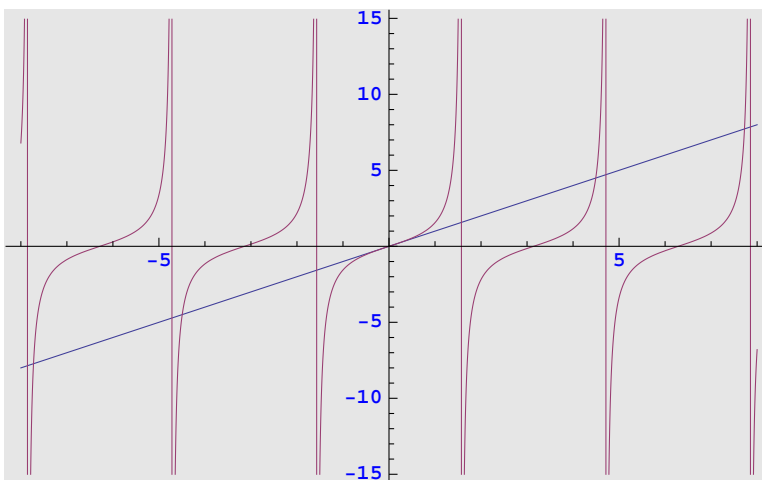


Figure 1: Graphs of  $\tan(x)$  and  $x$

To prove this analytically, choose any nonnegative integer  $k$  and consider the interval  $[(k + 1/2)\pi/L, (k + 3/2)\pi/L]$ . We know from properties of  $\tan$  that

$$\lim_{t \rightarrow \frac{\pi(k+1/2)}{L}^+} \tan(x) = -\infty, \quad \lim_{t \rightarrow \frac{\pi(k+3/2)}{L}^-} \tan(x) = \infty,$$

so for  $x$  slightly larger than  $(k+1/2)\pi/L$ ,  $\tan(x) < x/a$  and for  $x$  slightly smaller than  $(k+3/2)\pi/L$ ,  $\tan(x) > x/a$ . By the Intermediate Value Theorem, the functions must be share at least one point in the strip. Since the curves intersect in each of these strips, there are infinitely many intersections.

**Bonus:** Repeating some of part (a), we know we need to solve (1) for  $b$ , given a positive  $a$ , so we have

$$\tanh(2bL) = \frac{a}{b}.$$

These curves will always intersect as long as  $a > 0$ . Notice that  $\tanh(0) = 0$  and  $\lim_{b \rightarrow \infty} \tanh(2bL) = 1$ , whereas  $\lim_{b \rightarrow 0^+} a/b = \infty$  and  $\lim_{b \rightarrow \infty} a/b = 0$ , so the curves must cross. So the condition is that for any  $a > 0$ , there is a negative eigenvalue.

2. **(Strauss 4.3.2.)** Consider the eigenvalue problem with Robin boundary conditions

$$A'' + \lambda A = 0, \quad A'(0) - \alpha_0 A(0) = 0, \quad A'(L) + \alpha_L A(L) = 0.$$

(a) Show that zero is an eigenvalue if and only if  $\alpha_0 + \alpha_L = -\alpha_0 \alpha_L L$ .

(b) Compute the eigenfunction corresponding to this eigenvalue.

*Solution.*

(a) If we have a zero eigenvalue, then we have  $A'' = 0$  or

$$A(x) = \alpha x + \beta.$$

Plugging in the boundary conditions gives

$$\begin{aligned} \alpha - \alpha_0 \beta &= 0, \\ \alpha + \alpha_L(\alpha L + \beta) &= 0. \end{aligned}$$

We want to find a nontrivial solution for  $\alpha, \beta$  in this equation, or, as above, we need the matrix

$$\begin{pmatrix} 1 & -\alpha_0 \\ 1 + \alpha_L L & \alpha_L \end{pmatrix}$$

to have determinant zero, or

$$\alpha_L + \alpha_0(1 + \alpha_L L) = 0.$$

(b) If this determinant is zero, we know that the two equations we have are redundant, so we can solve either. The simpler to solve is the first, which gives  $\alpha = \alpha_0 \beta$ , and of course we will have one free choice for  $\beta$ . So one eigenfunction we can choose is

$$A(x) = \alpha_0 x + 1,$$

and we can of course choose any scalar multiple of this.

3. Solve the equation

$$\begin{aligned} u_t &= k u_{xx}, \quad x \in [0, \infty), \quad t > 0, \\ u(x, 0) &= \begin{cases} 1, & x \in (0, 1), \\ 0, & x > 1, \end{cases} \\ u(0, t) &= 0. \end{aligned}$$

*Solution.* We use the formula as derived in class (equation (6) in §3.1):

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_0^\infty \left( e^{-(x-y)^2/4kt} - e^{-(x+y)^2/4kt} \right) \phi(y) dy,$$

where  $\phi(y) = 1$  for  $y < 1$  and 0 otherwise. So we can also write

$$u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_0^1 \left( e^{-(x-y)^2/4kt} - e^{-(x+y)^2/4kt} \right) dy.$$

Let us consider the first term alone,

$$\frac{1}{\sqrt{4\pi kt}} \int_0^1 e^{-(x-y)^2/4kt} dy.$$

Using the change of variables  $s = (x - y)/\sqrt{4kt}$ , we obtain

$$\frac{1}{\sqrt{\pi}} \int_{\frac{x-1}{\sqrt{4kt}}}^{\frac{x}{\sqrt{4kt}}} e^{-s^2} ds,$$

or

$$\frac{1}{2} \left( \operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right) - \operatorname{erf} \left( \frac{x-1}{\sqrt{4kt}} \right) \right).$$

Now consider the second term,

$$\frac{1}{\sqrt{4\pi kt}} \int_0^1 e^{-(x+y)^2/4kt} dy.$$

Using the change of variables  $s = (x + y)/\sqrt{4kt}$ , we obtain

$$\frac{1}{\sqrt{\pi}} \int_{\frac{x}{\sqrt{4kt}}}^{\frac{x+1}{\sqrt{4kt}}} e^{-s^2} ds,$$

or

$$\frac{1}{2} \left( \operatorname{erf} \left( \frac{x+1}{\sqrt{4kt}} \right) - \operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right) \right).$$

Adding these together gives

$$-\frac{1}{2} \left( \operatorname{erf} \left( \frac{x+1}{\sqrt{4kt}} \right) - 2 \operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right) + \operatorname{erf} \left( \frac{x-1}{\sqrt{4kt}} \right) \right).$$

4. Consider the Schrödinger equation with Neumann boundary conditions:

$$iu_t = u_{xx}, \quad \frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(L, t) = 0.$$

Write out the general series solution for this equation as we have done for the heat and wave equations, i.e. separate variables, get ODEs in  $x$  and  $t$ , solve these problems, and take the linear combination.

*Solution.* We separate variables as usual, writing

$$u(x, t) = A(x)B(t).$$

Working this out gives us

$$A'' + \lambda A = 0, \quad A'(0) = A'(L) = 0, \quad B' + i\lambda B = 0.$$

for the  $A$  equation. We've solved this in class, and we know that the eigenvalues and eigenfunctions are

$$\lambda_n = \left( \frac{n\pi}{L} \right)^2, \quad A_n(x) = \cos \left( \frac{n\pi}{L} x \right), \quad n = 0, 1, 2, \dots$$

Solving the  $B$  equation gives

$$B_n(t) = B_n(0)e^{i\lambda_n t},$$

which oscillates. Notice of course that we can allow  $B_n(0)$  to be complex with no extra difficulties.

Thus we have

$$u_n(x, t) = C_n e^{i\lambda_n t} \cos\left(\frac{n\pi}{L}x\right),$$

and we form a general solution by linear combinations:

$$u(x, t) = \sum_{n=0}^{\infty} C_n e^{i\lambda_n t} \cos\left(\frac{n\pi}{L}x\right).$$